Photophysical Properties of Hexanuclear Copper(I) and Silver(I) Clusters

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The spectrophotometric properties, including emission spectra and lifetimes, absorption and diffuse reflectance spectra and excitation spectra, of the hexanuclear copper(I) and silver(I) clusters Cu₆(mtc)₂ (1), Ag₆(mtc)₂ (2), and Ag₆(dtc)₂ (3) (mtc = di-n-propylenethiocarbamate, dtc = di-n-propyldithiocarbamate) are reported. At room temperature I exhibits a strong single emission (e.g., in toluene Φₘₖ = 725 nm, τ = 1 μs), while at 77 K each of these is a bright emitter (spectral properties: I in toluene, Φₘₖ = 762 nm, τ = 20 μs; 2 in toluene, Φₘₖ = 607 nm, τ = 131 μs; 3 in toluene, Φₘₖ = 550 nm, τ = 21 μs). In analogy to previous assignments in tetranuclear Cu(I) clusters, the emission bands in the octahedral (M₄L₄) clusters are assigned to spin-forbidden transitions from triplet excited states which are mixed d-s and LMCT (ligand to metal charge transfer) in character modified by metal–metal interactions.

Introduction

During recent years there has been considerable interest in the photophysical properties of polynuclear d¹⁰ metal complexes, including those of copper(I), silver(I) and gold(I). One issue of special interest is the extent of metal–metal interaction in the ground and excited states. Although numerous studies have been concerned with the photophysical properties of polynuclear Cu(I) and Au(I) complexes, polynuclear complexes of the homologous Ag(I) metal center have received little attention, perhaps because of the sensitivity of silver compounds toward photo-decomposition. In addition, previous quantitative studies of polynuclear Cu(I) complexes have largely focused on compounds which contain a tetrahedral metal cluster. An extension of these investigations to hexanuclear octahedral metal clusters should be of considerable interest.

Described here are photophysical investigations of the hexaneric copper(I) and silver(I) complexes Ag₆(mtc)₂ (2), Ag₆(dtc)₂ (3), and Cu₆(mtc)₂ (1) (mtc = di-n-propylenethiocarbamate, dtc = di-n-propyldithiocarbamate, n-C₃H₇NCS₂), each of which has a structure analogous to A shown in Figure 1.

Experimental Section

Materials.

All chemicals used in syntheses were reagent grade. Emission, excitation and lifetime data were recorded under a dinitrogen atmosphere. Solvents for spectroscopic studies were redistilled before use.

Compounds 1 and 2 were synthesized and purified according to published methods. Compound 3 was synthesized by Akerström’s method and purified by column chromatography (mobile phase toluene/diethyl ether (1:1); stationary phase Sigma No. A-9003 Chromatographic Alumina Neutral Type WN-3 activity grade 1). Exposure to light during synthesis and purification of 2 and 3 was avoided.

Melting points and crystal lattice constants of 1 and 2 were in good correspondence with data published by Akerström,° Hesse,‡ and Jennische.‡ (1) published, mp 180–183 °C, a = 13.425 Å, b = 19.357 Å, c = 11.630 Å, α = γ = 90.00°, β = 91.72°; 1 experimental, mp 180–182 °C, a = 13.381 Å, b = 19.390 Å, c = 11.615 Å, α = γ = 90.00°, β = 90.21°; 2 published, mp 133–134 °C, a = 13.313 Å, b = 19.904 Å, c = 11.901 Å, α = γ = 90.00°, β = 93.28°; 2 experimental, mp 129–131 °C, a = 13.241 Å, b = 19.672 Å, c = 11.766 Å, α = γ = 90.00°, β = 90.33°. The melting point of 3 (mp 152–153 °C) was in agreement with published data (mp 154–155 °C). The elemental analyses of 1–3 are in good agreement with calculated data. Anal. Caled for 1, C₆H₆N₂O₂S₅Cu₂: C, 37.57; H, 6.31; N, 6.26. Found: C, 37.61; H, 6.21; N, 6.29. Caled for 2, C₆H₆N₂O₂S₅Ag₂: C, 31.36; H, 5.26; N, 5.26. Found: C, 31.43; H, 5.23; N, 5.26. Caled for 3, C₆H₆N₂S₂Ag₂: C, 25.99; H, 4.97; N, 4.53. Found: C, 25.92; H, 4.80; N, 4.82.

Solutions of 1 and 2 were thermally stable but underwent slow deterioration when irradiated with white light. Solid samples of 1 were sensitive to irradiation with a Nd:YAG laser (355 nm). Solutions and solid samples of 3 were highly sensitive to irradiation with white light or a Nd:YAG laser (355 nm). Solutions of 3 were much less stable than those of 1 or 2. Pure 3 exhibited a green emission (about 550 nm) when irradiated with UV light at 77 K. The decomposed product exhibited a very broad orange to red emission (up to 700 nm) under the same conditions.

Instrumentation.

Absorption spectra of solutions were collected with a HP 8452 diode array UV/vis spectrophotometer in 1-cm quartz cells. Diffuse reflectance (DR) spectra of solids were recorded on a Cary 14 UV/vis spectrophotometer equipped with an integrating sphere and On-Line Instrument Systems (OLIS) computer control system in the laboratory of G. Stucky at UCSB. Solid MgCO₃ was used as reference, and samples were diluted with MgCO₃ (∼5–10% sample by weight).

Emission spectra were collected with a Spex Fluorolog 2 spectrofluorimeter (150-W Xe lamp, two single monochromators, water-cooled Hamamatsu R928A photomultiplier (PMT) configured for photon counting) with a Spex Datamate II data station. Emission was passed through a Corning glass cutoff filter (3-72 or 3-71) prior to the PMT. Emission spectra were normalized but were not corrected for PMT response. Quartz or Pyrex tubes (5 mm) sealed with septa or flame-sealed were used as sample cells. The range of solution concentrations used was 10⁻⁴–10⁻⁵ M. Data at 77 K were recorded by cooling the sample with liquid nitrogen in a Pyrex Dewar flask. Excitation spectra were collected with the Spex Fluorolog 2 system and were normalized and corrected for lamp intensity variation by the ratio method with Rhodamine 6G reference.

Lifetime data were collected with a Quantas Ray DCR-1A Q-switched Nd:YAG pulsed laser (operating at 10 Hz) as the excitation source. The power output at 355 nm was 40–50 mJ/pulse for samples of 1 and <10 mJ/pulse for samples of 2 and 3. A HG-1 harmonic generator and a PH1-1 harmonic separator were used to isolate the desired third harmonic (355 nm). Specific wavelength dichroic mirrors mounted on beam steering towers were used to direct the excitation beam and to ensure laser-quality at the sample. The sample emission was monitored at a right angle to the excitation source. The sample emission was filtered with appropriate filters (Schott plastic filter KV-388 and Corning glass cutoff filters 3-72 and 3-69, Schott glass filters LP 40 and LP 34) to reduce laser scatter prior to entering a Spex Model 1680 Doublet grating monochromator bleached at 500 nm. The emission intensity was monitored at selected wavelengths with an RCA 8852 fast-response PMT. The PMT output was terminated (50 Ω) into a Tektronix 7912AD


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transient digitizer (1 and 3) or a Le Croy 9400 Dual 125-MHz digital oscilloscope (2). The Tektronix 7912AD transient digitizer was equipped with a 7890P programmable time base and a TA13 differential comparator amplifier. The signal-averaged data (64 shots per 512 point array) for 1 and 16 shots for 3) were processed using a Zenith ZF-158-42 PC. Data were analyzed by exponential curve fitting. Data for 2 (16 shots per 512 point array) were evaluated manually by a plot of ln (intensity) versus time. Samples for lifetime measurements were prepared in the same manner as those for emission spectra. Unless otherwise noted, lifetime measurements were reproducible to within ±10%.

High-resolution powder diffraction data were collected for each sample with a Scintag PADX automated powder X-ray diffractometer (Model XPH-105) using Cu Kα radiation.

Results

Absorption and Diffuse Reflectance Spectra. The absorption spectra of solution samples of 1–3 and the diffuse reflectance spectra of the solids of the same materials are summarized in Table I. Examples of the solution spectra are shown in Figure 2. These compounds display strong absorptions in the UV region, which appear to be ligand-localized transitions, shifted to somewhat lower energy, plus much weaker bands or shoulders at longer wavelengths. The spectral patterns for the two monothiocarbamate clusters 1 and 2 were similar, with the longer wavelength absorption appearing as a peak at ~340 nm with a shoulder at ~380 nm for Cu_mtc and as a peak at ~310 nm and shoulder at ~360 nm for Ag_mtc. These bands were relatively insensitive to solvent but did shift modestly to higher energy in the more polar media. Diffuse reflectance spectra of the respective solids demonstrated the transitions at nearly the same wavelengths, although the shoulders appeared in the DR spectra as peaks for both 1 and 2 (Figure 3).

The solution absorption spectra of the dithiocarbamate complex Ag_dtc, in different solvents show only a broad shoulder at ~336 nm in the near-UV region (Figure 2), although the absorption tails to much longer wavelength. However, this band appears clearly as a peak at 339 nm in the DR spectrum of solid 3 (Figure 3). Attempts to record the spectrum of 3 in acetonitrile or ethanol were unsuccessful owing to low solubility.

Emission Spectra and Lifetimes. Emission spectra and lifetime data are summarized in Tables II and III. The hexanuclear copper(I) cluster 1 exhibits a bright red emission at room temperature, both in toluene solution ($\lambda_{\text{max}} = 725$ nm) and as a solid ($\lambda_{\text{max}} = 706$ nm). The toluene solution-phase emission spectra of 1 at 294 and 77 K are shown in Figure 4. Lowering the temperature leads to a narrowing of the emission band, as is typically seen ($\Delta \lambda_{1/2} = \text{full width at half-height} = 2.7 \times 10^3$ cm$^{-1}$ at 294 K and 1.7 $\times 10^3$ cm$^{-1}$ at 77 K); however, this is accompanied by an atypically red shift of $\lambda_{\text{max}}$ to 762 nm. Both at room temperature and at 77 K the emission spectrum proved to be independent of the excitation wavelength $\lambda_{\text{exc}}$. Exponential decay was observed at both temperatures with the respective measured lifetimes in toluene being 0.97 and 19.9 μs. The emission properties of 1 in ethanol solution and as a solid were comparable (Tables II and III).

The hexanuclear silver(I) complexes 2 and 3 were found to emit only at 77 K. The bright, unstructured orange emission from the
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**Table II.** Emission and Excitation Data for Cu₆(mtc)₆ (1), Ag₆(mtc)₆ (2), and Ag₆(dtc)₆ (3)

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<th>compd</th>
<th>T°</th>
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<th>λₑₓᵖ max</th>
<th>Δνₑₓᵖ</th>
<th>λₑₓᵖ sh</th>
<th>ΔE°</th>
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<td>1.7</td>
<td>357 (p)</td>
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<td>toluene</td>
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<td>2.0</td>
<td>353 (sh)</td>
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* Temperature in K. ¹Uncorrected) emission maximum in nm (excitation wavelength = λₑₓᵖ). ²Full width at half-height of the emission band in 10⁴ cm⁻¹. ³Excitation maxima in nm monitored at λₑₓᵖ; p = peak, sh = shoulder. ⁴Energy difference (in 10⁴ cm⁻¹) between emission maxima and lowest excitation maximum.

**Table III.** Emission Lifetimes of Cu₆(mtc)₆ (1), Ag₆(mtc)₆ (2), and Ag₆(dtc)₆ (3)

<table>
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<th>compd</th>
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<th>τ¹</th>
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<td>77</td>
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<td>77</td>
<td>solid</td>
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<td></td>
<td>toluene</td>
<td>156</td>
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<td>Ag₆(dtc)₆</td>
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<td></td>
<td></td>
<td>EtOH/CH₂Cl₂</td>
<td>12.9</td>
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</table>

* Temperature in K. ¹τ in μs; monitored at λₑₓᵖ max, unless noted; experimental uncertainties were ±10%. ²Biexponential decay.

Figure 5. (a) Emission spectrum of Ag₆(mtc)₆ (2) in toluene solution at 77 K (λₑₓᵖ = 360 nm). (b) Excitation spectrum of 2 in toluene at 77 K, monitored at λₑₓᵖ max. Intensities were normalized for comparison.

Figure 6. (a) Emission spectrum of the dithiocarbamate cluster Ag₆(dtc)₆ (3) in 77 K toluene (λₑₓᵖ = 430 nm). (b) Excitation spectrum of 3 in toluene at 77 K, monitored at λₑₓᵖ max. Intensities were normalized for comparison.

same conditions. The bright green emission from the dithiocarbamate cluster (Ag₆(dtc)₆) occurs at even higher energy with λₑₓᵖ max values of 545 and 550 nm (Figure 6) as a solid and in a toluene solution, respectively. For both clusters, the emission spectra were independent of the excitation wavelength and showed only modest shifts to a shorter λₑₓᵖ max when ethanol was used as the solvent instead of toluene. Lifetime measurements showed the emission from solid 2 to have a lifetime of ~100 μs but to be nonexponential, while single-exponential decays were observed in ethanol and toluene solutions. Notably, these were nearly an order of magnitude longer (131 ± 10 and 156 ± 12 μs, respectively) than the measured lifetimes of the Cu(I) analogue 1 under similar conditions (see above) or of the dithiocarbamate complex 3, which gave lifetimes of 8.4 μs (solid), 21.3 μs (toluene solution), and 12.9 μs (ethanol/dichloromethane solution). Bandwidths and other relevant photophysical data are summarized in Table II.

Excitation Spectra. The excitation spectra for the emissions from toluene solutions of 1-3 (monitoring at λₑₓᵖ max in each case) are displayed in Figures 4-6, and the λₑₓᵖ max values determined in other media are summarized in Table II. In each case, the excitation spectra displayed two maxima. For the Cu₆(mtc)₆ in 77 K toluene, the λₑₓᵖ max values were 357 and 430 nm, shifted only slightly from those seen in 294 K solution (351 and 425 nm (sh)), and these match approximately the λₑₓᵖ values (342 and 430 nm (sh)) seen in the 294 K toluene solution absorption spectra of 1. Similarly, the excitation maxima for Ag₆(mtc)₆ in 77 K toluene solution occur at wavelengths (325 and 360 nm) approximating the λₑₓᵖ max of 294 K toluene solutions (312 and 360 nm). However, the same cannot be claimed for the 77 K toluene solution of the (dithiocarbamate)silver(I) cluster, the excitation spectrum of which displayed a λₑₓᵖ max at 431 nm, a wavelength for which neither the solution absorption spectra nor the solid diffuse reflectance spectrum (which generally is more sensitive to low energy, forbidden transitions) indicated a transition. The excitation spectra of 3 did indicate a weaker λₑₓᵖ max at ~350 nm, which does correspond approximately to the lowest energy shoulder in the ab-
The Stokes shift from the lowest energy excitation band maximum to the emission band maximum gives a $\Delta E$ of about $10^{-4}$ cm$^{-1}$ for the two monothiocarbamate clusters, but the higher energy of the emission and the lower energy excitation maximum makes the Stokes shift for the dithiocarbamate complex 3 much smaller.

**Discussion**

The following features should be considered in examining the photophysical properties of the (M)$\delta$ clusters. First, for each complex, the absorption spectrum shows a lowest energy shoulder or peak in the near-UV region at wavelengths longer than those of the free-ligand absorption bands.

These transitions, which are also prominent in the diffuse reflectance spectra of the solids, follow the energy order $\text{Ag}_2(dtc)_2 > \text{Ag}_2(mtc)_2 > \text{Cu}_2(mtc)_2$. Second, at 77 K each of the clusters displays an unstructured, visible-region emission band, the energy order again following the sequence $\text{Ag}_2(dtc)_2 > \text{Ag}_2(mtc)_2 > \text{Cu}_2(mtc)_2$. The position of this band is relatively independent of the solvent medium. The Cu$_{\delta}$ cluster displays this emission at ambient temperature as well, in contrast to the Ag$_{\delta}$ clusters. Third, Stokes shifts from the excitation maxima, especially for the two mtc$^-$ complexes, are quite large and suggest significant structural distortion between the emitting and ground states. Fourth, luminescence lifetimes at 77 K are relatively long ($>1 \mu$s), implying that the transitions involved are forbidden, most likely the result of emissions from triplet excited states.

Structurally, the metal atoms in the hexameric clusters 1 and 2 form nearly regular octahedra with metal-metal distances (1.2.70-3.06 Å; 2.94-3.28 Å) somewhat larger than in the bulk metals (Cu, 2.56 Å; Ag, 2.88 Å). If the comparison is made to van der Waals radii of the metals (1.40 and 1.70 Å, respectively), the distances are closer; indeed, the average Ag–Ag distance in 2 is somewhat less than the summed van der Waals radii (3.4 Å). For 3, the (Ag$^+$)$_2$ core is more distorted from a regular octahedron, with a set of six short (2.91–3.20 Å) and six long (3.85–4.01 Å) Ag–Ag edges. Six of the eight triangular faces of the (M)$^+$ octahedral cores are capped by the bidentate ligands with a sulfur atom bridging one M–M edge and the oxygen of mtc$^-$ or the second sulfur of dtc$^-$ terminally bound to the third M on the face, as illustrated in Figure 1. Cryoscopic and ebulliometric measurements have verified that the hexameric units of the three compounds are preserved in solution.

This conclusion is supported by the similarities observed here for the emission spectra of each of these cluster compounds in the solid and in solution. Notably, there is a strong qualitative similarity between the absorption and emission spectral features of 1–3, which suggests that these may be characteristic properties of these (M)$\delta$ clusters.

A detailed interpretation of the complicated electronic structures of these clusters will not be attempted here. However, a qualitative model which accounts for some features in the absorption and emission spectral properties will be developed in the following discussion.

Possible assignments for the excited states relevant to the emissions observed are those originating from ligand-centered $\sigma$–$\pi^*$ (LC) transitions, from metal–ligand charge-transfer (LMCT) transitions, or from either metal-to-ligand (MLCT) or ligand-to-metal charge-transfer (LMCT) transitions. Of these, neither the LC nor the MLCT assignment would appear likely. The LC transitions appear to be too high in energy, especially for the mtc$^-$ ligand ($\lambda_{\text{abs}} > 200$ nm), to be responsible for the lowest energy absorptions or the emission bands. The LC absorption bands of other d10 complexes, e.g., Zn(dtc)$_2$, have energies close to those of the free anion; furthermore, it is notable that the free ligands as sodium salts do not display any visible-range emissions. MLCT excited states have been identified as important in the photophysical properties of the copper(I) complexes of easily reducible ligands such as phenanthroline, but such an assignment does not seem reasonable for the mono- and dithiolato ligands, especially in light of the relevant bands for the Cu(I) and Ag(I) complexes having similar energies. Generally, Cu(I) is much more easily oxidized, given the higher energy of the Cu 3d orbitals relative to those of the Ag 4d orbitals. Furthermore, the s* orbitals of mtc$^-$ and dtc$^-$ are too high in energy to be realistically expected to act as acceptor orbitals for low-lying MLCT transitions.

The above considerations leave the LMCT and metal-centered $d \rightarrow s$ transitions as the most likely assignments. Since the electron affinities of the free Cu$^+$ and Ag$^+$ ions are comparable, to a first approximation LMCT transitions of the Cu(I) and Ag(I) complexes of similar ligands should have comparable energies, although, for the hexanuclear clusters, the molecular orbitals would certainly be dependent on the extent of mixing due to metal–metal and metal–ligand covalent interactions. The optical electronegativities of dithiocarbamate ligands are relatively small (2.7), just higher than that of iodide (2.5). Given that the thiocarbamate complexes of more oxidizing metal ions such as Cu(II) or Ag(II) are strongly colored owing to LMCT absorption bands in the visible region, it is reasonable to speculate that the analogous bands for the Cu(I) and Ag(I) clusters would appear in the near-ultraviolet region of the spectrum.

Alternatively, since $d \rightarrow s$ transitions of the free ions Cu$^+(g)$ and Ag$^+(g)$ occur at relatively low energies (21.93 $\times 10^3$ and 26.26 $\times 10^3$ cm$^{-1}$, respectively, for the lowest energy spin-forbidden and spin-allowed transitions of the former, 39.16 $\times 10^3$ and 46.05 $\times 10^3$ cm$^{-1}$ for the analogous transitions of the latter), $d \rightarrow s$ states would also appear to be reasonable candidates for low-lying excited states among the (M)$\delta$ clusters. The appearance of both the lowest wavelength absorption band and the emission band for the Cu(I)–mtc cluster 1 at a somewhat lower energy than for $d \rightarrow s$ transitions (Table I). Notably, the Stokes shift from the lowest energy excitation band maximum to the emission band maximum gives a $\Delta E$ of about $10^{-4}$ cm$^{-1}$ for the two monothiocarbamate clusters, but the higher energy of the emission and the lower energy excitation maximum makes the Stokes shift for the dithiocarbamate complex 3 much smaller.

**Figure 7.** Metal orbitals of the M$_6$ octahedron (diagram adapted from ref 30).

\[ \text{splitting of frontier metal orbitals in octahedral } M_6 \text{ cluster} \]
the Ag(I) analogue 2 would be consistent with the ordering of the $d \rightarrow s$ energies of the mononuclear free ions (above), but the differences are much smaller. Ligand coordination would be expected to raise $s$-orbital energies substantially, but the effect of metal–metal interactions in the $(M^I)_2$ octahedra would lead to substantial splitting of the $d$ and (especially) $s$ orbitals in a manner which may lead to marked effects on the expected molecular orbital energies. A diagram representing the resulting cluster metal MO’s is displayed in Figure 7 and illustrates the stabilization of the lowest energy unfilled $s$ orbital and corresponding destabilization of the highest energy filled $d$ orbital. It is likely that the metal–metal interactions in the silver clusters are substantially greater than in the respective copper cluster because, relative to the sums of their respective van der Waals radii, the Ag–Ag distances, especially in 2 but also in 3, are significantly shorter than the Cu–Cu distance in 1. This enhanced interaction should decrease the $d \rightarrow s$ separation accordingly.

Earlier studies from these two laboratories$^{23}$ have identified strong, low-energy bands in the luminescence spectra of the Cu(I) cluster Cu$_4$I$_4$(py)$_4$ (4) and a series of related Cu$_4$L$_4$ analogues where $L$ is either a saturated or aromatic amine. Recent ab initio calculations$^{26}$ have concluded that the emitting excited state can be viewed to be of mixed character consisting of roughly equal parts of iodide to copper charge transfer (XMCT, $X = I$) and metal-centered $d \rightarrow s$ contributions, the energies markedly influenced by the metal–metal interactions in the excited state. Blasse$^{27}$ has reached a similar conclusion. The idea of mixed $d \rightarrow s$/XMCT character of electronic transitions was advanced some years ago by Jørgensen$^{28}$ to explain the absorption spectra of Cu(I) and Ag(I) halide salts. For example, when Ag$^+$ is doped into alkali metal halides ($X^-$), the longest wavelength $d \rightarrow s$ transition follows the energy order $Cl^- > Br^- > I^-$. The iodide salts show the largest shift to lower energy due to greater mixing of the metal $d$ orbitals with the halide $p$ orbitals and, hence, a greater LMCT contribution to the transition.

Also relevant to the present discussion is a self-consistent charge and configuration MO investigation$^{29}$ of Cu(I) clusters of the type having a Cu$_4$S$_{12}$ $^{4-}$ core with the eight copppers in a cubane configuration with Cu–Cu distances of 2.83 Å and the sulfurs coming from bidentate ions such as 1,1-dicyanoethylene-2,2-dithiolate. This study concluded that the HOMO’s are largely sulfur in character while the LUMO is primarily copper 4$s$ and 4$p$ and sulfur 3$p$ orbitals. These results would imply that the HOMO’s for the octahedral clusters described in the present case are also largely ligand in character. However, the ab initio calculations described for the Cu(I) tetrahedra$^{26}$ clearly indicate considerable ground-state charge delocalization from ligand anions to the metal core and much less than a full electron transfer from ligand to metal in the lowest energy excited state. Thus, in such a case, viewing the ES conveniently as simply the product of a one-electron HOMO $\rightarrow$ LUMO transition may present an incomplete picture of the changes wrought by the electronic excitation.

In the context of the calculations described above for Cu$_4$I$_4$(py)$_4$, a "cluster-centered" (CC) state with both $d \rightarrow s$ and LMCT character would appear to be an accurate, although qualitative, description of the luminescent excited states of 1–3. The $a_g$ lowest unoccupied MO (MO) of the $(M^I)_2$ octahedra (Figure 7) would be largely metal $s$ orbital in character and would be a bonding MO with respect to the metal–metal interactions. Therefore, a $d \rightarrow s$ transition, a LMCT transition, or the presumed CC transition of mixed character would lead to population of this orbital with a resulting contraction of the $(M^I)_2$ octahedron. Such a distortion would be one factor accounting for the large Stokes shift from the excitation maxima to the emission maxima of the monothiocarbarate complexes 1 and 2. We do not have any explanation for the smaller, but still substantial, Stokes shift for 3, although it should be noted that, in the ground state for this species, the $(M^I)_4$ core is already significantly distorted from the nearly regular octahedra which characterize the other two. The relatively long emission lifetimes at 77 K for 1–3 (Table III) are consistent with those expected for a spin-forbidden transition in such metal complexes. The emitting state is thus concluded to be a triplet CC ES of mixed $d \rightarrow s$ and LMCT character.

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Raman spectra of molten Hg$_2$X$_2$–Hg (X = Cl, Br, I) systems have been obtained at compositions up to 30 mol % in Hg from 550 to 818 K. The dissolution of mercury in mercury halides gives rise to resonance-enhanced Raman bands which were interpreted for Hg$_2$X$_2$ type molecular species formed in all mercury compositions and Hg$_2$X$_2$ type molecules formed at high mercury mole fractions. Spectra were also obtained from HgX$_2$–Hg$_2$X$_2$ (X = F, Cl, Br, I) melts and were attributed to mixed mercury(I) (sub)halide molecules Hg$_2$XX’ formed in the melt. The Hg$_2$X$_2$ and Hg$_2$XX’ molecules possess a linear symmetry, and the Hg–Hg stretching frequencies for all 10 molecules were found to be between $\sim$180 cm$^{-1}$ (Hg$_2$F$_2$) and $\sim$100 cm$^{-1}$ (Hg$_2$I$_2$). A linear Hg$_2$ chain is formed in the Hg$_2$X$_2$ molecules bound to two terminal halides. The formation of Hg$_2$ chains was further confirmed by the Raman spectra of Hg$_2$(ACl)$_2$ melts. In the suggested state, it is suggested that in the melt mixtures intermolecular interactions between Hg$_2$X$_2$ and Hg$_2$X$_4$ molecules lead to an alteration of oxidation states which account for a "hopping" like conduction.

**Introduction**

At elevated temperatures many metals are to some extent soluble in their molten halide salts, forming, in most cases, colored solutions. The basic physical properties and the equilibrium phase diagrams of a large number of metal–metal halide (M–MX$_n$) systems were determined in the 1960s, and most of this work was summarized in reviews by Bredig$^1$ and Corbett.$^2$ Some interesting...