

METHODS FOR THE CONSTRUCTION OF EXTRAPOLATION PROCESSES

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- General Approach
- Iterative Approach
- Variational Approach
- Perturbative Approach

MANY THANKS TO

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Auswertiges Amt

Development of Methods:

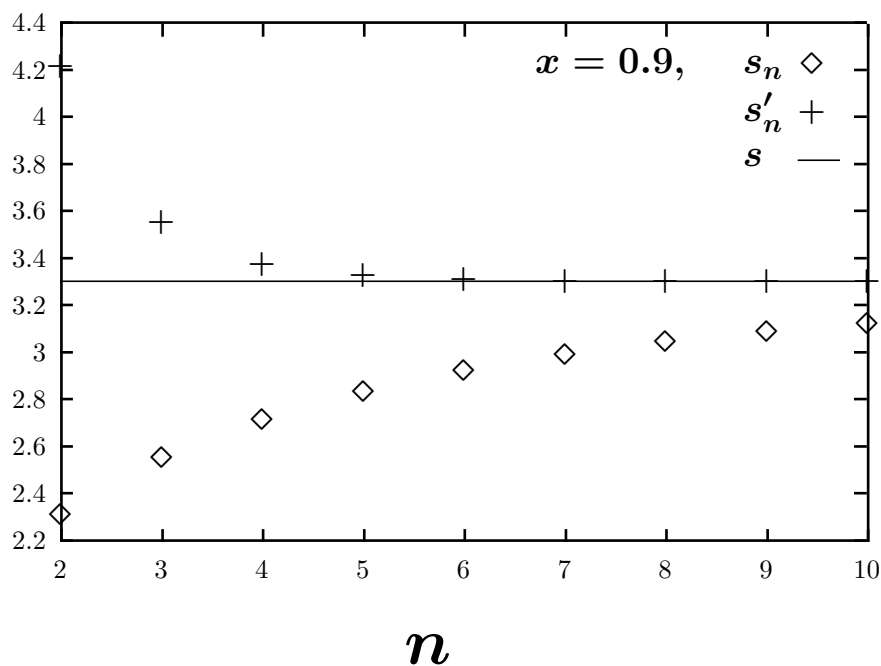
Extrapolation Methods

Problem:
Slow Convergence

Example

$$s_n = 1 + \frac{x}{1} + \frac{x^2}{2} + \frac{x^3}{3} + \dots + \frac{x^n}{n}$$

$$s = \lim_{n \rightarrow \infty} s_n = 1 - \ln(1 - x)$$



Sequence transformation $\{s'_n\} = T(\{s_n\})$:
Accelerate convergence

Development of Methods: *Extrapolation Methods*

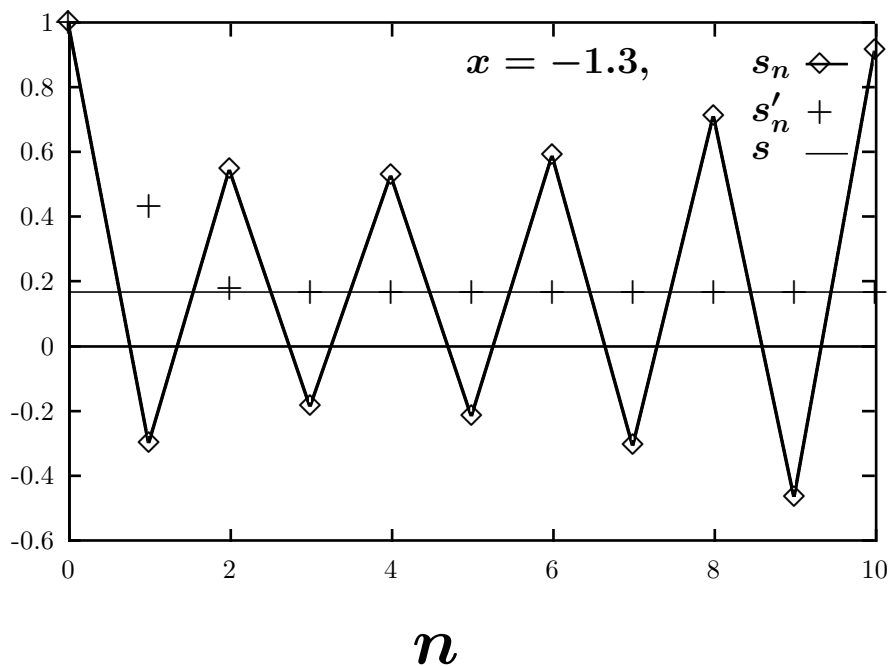
Problem:

No Convergence

Example:

$$s_n = 1 + \frac{x}{1} + \frac{x^2}{2} + \frac{x^3}{3} + \dots + \frac{x^n}{n}$$

$$s = \lim_{n \rightarrow \infty} s_n = 1 - \ln(1 - x)$$



Sequence transformation $\{s'_n\} = T(\{s_n\})$:
Achieve convergence

PROBLEMS

- Extrapolation
- Convergence acceleration
- Summation of divergent series
- No universal method available
- Linear methods often not efficient enough
- Accuracy of single method unknown in practice

AIMS

- problem adapted
- nonlinear
- efficient

APPROACH

- Model sequences
- Iteration of simple methods
- Comparison of results of various methods

BASIC PRINCIPLE

EXAMPLE

Model sequences

$$\sigma_n = \sigma + \mathcal{M}_n(\vec{c}, \vec{p}(n))$$

$$\sigma_n = \sigma + c\omega_n$$

Exact Calculation of σ

$$\sigma = T_n(\{\sigma_n\}, \vec{p}(n))$$

$$\sigma = \sigma_{n+1} - \omega_{n+1} \frac{\sigma_{n+1} - \sigma_n}{\omega_{n+1} - \omega_n}$$

Sequence transformation for problem $\{s_n\}$

$$s'_n = T_n(\{s_n\}, \vec{p}(n))$$

$$s'_n = s_{n+1} - \omega_{n+1} \frac{s_{n+1} - s_n}{\omega_{n+1} - \omega_n}$$

Acceleration for

$$\{s_n\} \approx \{\sigma_n\}$$

$$\frac{s_n - s}{\omega_n} = O(1)$$

ANNIHILATION OPERATORS

$$\sigma_n = \sigma + \omega_n \sum_{j=0}^k c_j \phi_j(n)$$

with linear operator

$$\mathcal{O}_n(\phi_j(n)) = 0, \quad j = 0, \dots, k$$

then

$$\mathcal{O}_n \left(\frac{\sigma_n - \sigma}{\omega_n} \right) = 0$$

or

$$\sigma = \frac{\mathcal{O}_n(\sigma_n/\omega_n)}{\mathcal{O}_n(1/\omega_n)} \implies s'_n = \frac{\mathcal{O}_n(s_n/\omega_n)}{\mathcal{O}_n(1/\omega_n)}$$

CONVERGENCE CLASSIFICATION

$$\lim_{n \rightarrow \infty} \frac{s_{n+1} - s}{s_n - s} = \rho$$

linearly convergent: $0 < |\rho| < 1$

logarithmically convergent: $\rho = 1$

REMAINDER ESTIMATES

$$\frac{s_n - s}{\omega_n} = O(1), \quad n \rightarrow \infty$$

Series

$$s_n = \sum_{j=0}^n a_j \rightarrow s, \quad n \rightarrow \infty$$

Levin

$$\begin{aligned} t_{\omega_n} &= a_n, \\ u_{\omega_n} &= (n + \beta)a_n, \quad \beta > 0, \\ v_{\omega_n} &= \frac{a_n a_{n+1}}{a_n - a_{n+1}}. \end{aligned}$$

Smith and Ford

$$\tilde{t}_{\omega_n} = a_{n+1}$$

Asymptotically related series

$$\hat{a}_n \sim a_n$$

$$\hat{s}_n = \sum_{j=0}^n \hat{a}_j \rightarrow \hat{s}, \quad n \rightarrow \infty$$

Linear remainder estimates

$${}^{lt}\omega_n = \hat{a}_n,$$

$${}^{lu}\omega_n = (n + \beta)\hat{a}_n, \quad \beta > 0,$$

$${}^{lv}\omega_n = \frac{\hat{a}_n \hat{a}_{n+1}}{\hat{a}_n - \hat{a}_{n+1}},$$

$${}^{l\tilde{t}}\omega_n = \hat{a}_{n+1}$$

Kummer-type remainder estimate

$${}^k\omega_n = \hat{s}_n - \hat{s}$$

Tails

$$T_n = a_{n+1} + a_{n+2} + \cdots = s - s_n$$

$${}^{\tilde{t}}T_n = s_{n+1} - s_n = a_{n+1} + 0 + \cdots,$$

$${}^{l\tilde{t}}T_n = \hat{s}_{n+1} - \hat{s}_n = \hat{a}_{n+1} + 0 + \cdots,$$

$${}^kT_n = \hat{s} - \hat{s}_n = \hat{a}_{n+1} + \hat{a}_{n+2} + \cdots.$$

Example

$$F_m(z) = \sum_{j=0}^{\infty} (-z)^j / j!(2m + 2j + 1) ,$$

$${}^k\omega_n = \left((1 - e^{-z})/z - \sum_{j=0}^n (-z)^j / (j + 1)! \right) .$$

n	s_n	${}^u\omega_n$	${}^t\omega_n$	${}^k\omega_n$
5	-13.3	0.3120747	0.3143352	0.3132981
6	14.7	0.3132882	0.3131147	0.3133070
7	-13.1	0.3132779	0.3133356	0.3133087
8	11.4	0.3133089	0.3133054	0.3133087
9	-8.0	0.3133083	0.3133090	0.3133087

$$z = 8 , m = 0$$

${}_2J$ -Transformation

ITERATIVE SEQUENCE TRANSFORMATIONS

Idea:

- Simple sequence transformation

$$\{s'_n\} = T(\{s_n\})$$

- Iteration:

$$\{s''_n\} = T'(\{s'_n\}), \{s'''_n\} = T''(\{s''_n\}), \dots$$

- Example: Aitken Δ^2 method

$$s_n^{(1)} = s_n - \frac{(s_{n+1} - s_n)^2}{s_{n+2} - 2s_{n+1} + s_n}.$$

Iteration

$$\mathcal{A}_0^{(n)} = s_n,$$

$$\mathcal{A}_{k+1}^{(n)} = \mathcal{A}_k^{(n)} - \frac{(\mathcal{A}_k^{(n+1)} - \mathcal{A}_k^{(n)})^2}{\mathcal{A}_k^{(n+2)} - 2\mathcal{A}_k^{(n+1)} + \mathcal{A}_k^{(n)}}.$$

Problem:

- Iteration not unique, for example

$$s_n^{(k+1)} = T^{(k)}(\{s_n^{(k)}\})$$

with $T^{(0)} = T$.

- Example

$$\begin{aligned}\overline{\mathcal{A}}_0^{(n)} &= s_n, \\ \overline{\mathcal{A}}_{k+1}^{(n)} &= \overline{\mathcal{A}}_k^{(n)}\end{aligned}$$

$$-X_k \frac{(\overline{\mathcal{A}}_k^{(n+1)} - \overline{\mathcal{A}}_k^{(n)})^2}{\overline{\mathcal{A}}_k^{(n+2)} - 2\overline{\mathcal{A}}_k^{(n+1)} + \overline{\mathcal{A}}_k^{(n)}}$$

with $X_0 = 1$ “allowed”,
i.e., $X_k = (2k + 1)/(k + 1)$.

HIERARCHICAL CONSISTENCY

Example:

- Simple transformation:

$$s'_n = s_{n+1} - \omega_{n+1} \frac{s_{n+1} - s_n}{\omega_{n+1} - \omega_n}$$

- Problem: ω'_n ?

- For

$$\sigma_n = \sigma + \omega_n(c_0 + c_1 r_n)$$

we have

$$\sigma'_n = \sigma - c_1 \frac{\omega_n \omega_{n+1}}{\omega_{n+1} - \omega_n} (r_{n+1} - r_n)$$

- Hence

$$\omega'_n = -\frac{\omega_n \omega_{n+1}}{\omega_{n+1} - \omega_n} (r_{n+1} - r_n)$$

• Iteration $\implies \mathcal{J}$ transformation

$$\begin{aligned}
 s_n^{(0)} &= s_n, & \omega_n^{(0)} &= \omega_n, \\
 s_n^{(k+1)} &= s_n^{(k)} - \omega_n^{(k)} \frac{s_{n+1}^{(k)} - s_n^{(k)}}{\omega_{n+1}^{(k)} - \omega_n^{(k)}}, \\
 \omega_n^{(k+1)} &= -\frac{\omega_n^{(k)} \omega_{n+1}^{(k)}}{\omega_{n+1}^{(k)} - \omega_n^{(k)}} \delta_n^{(k)}, \\
 \mathcal{J}_n^{(k)}(\{s_n\}, \{\omega_n\}, \{r_n^{(k)}\}) &= s_n^{(k)}
 \end{aligned}$$

- for power series,
- very versatile and powerful.
- related to E algorithm with model

$$\sigma_n = \sigma + \sum_{j=0}^k c_j g_j(n)$$

- Analysed in
H. H. H. Homeier, A hierarchically consistent, iterative sequence transformation, Numer. Algo. 8 (1994) 47-81.
—, Analytical and numerical studies of the convergence behavior of the \mathcal{J} transformation, J. Comput. Appl. Math. 69 (1996) 81-112.
—, Determinantal representations for the \mathcal{J} transformation, Numer. Math. 71 (1995) 275-288.

Numerical example:

$$G_n = \int_0^1 \left\{ \int_{-\infty}^{\infty} \frac{\Gamma/\pi}{(x - \Delta\omega)^2 + \Gamma^2} \left(\exp \left[-\frac{Y q^2 \Gamma/\pi}{x^2 + \Gamma^2} \right] - 1 \right) dx \right\} q^n dq$$

Taylor series in $Y \implies$ Power series
Acceleration with \mathcal{J} transformation

n	s_n	$s_0^{(n)}$
15	-3047434.	-0.16361565
16	5412146.	-0.16361782
17	-9099655.	-0.16361716
18	14525645.	-0.16361732
19	-22070655.	-0.16361728
20	31994427.	-0.16361729
∞	-0.16361729	-0.16361729

$n = 2, \Delta\omega = 5, \Gamma = 1, Y = 100, r_k^{(k)} = 1/(n + 1 + k), k$ variant

More general:

- Hierarchy of model sequences: The higher, the more parameters
- Iteration of simple transformation T is consistent, if
 - T is exact for lowest level sequences
 - Variant of T maps higher to lower level

EXAMPLE:

Hierarchy for \mathcal{J} transformation

$$\begin{aligned}\sigma_n = & \sigma + \omega_n \left(c_0 + c_1 \sum_{n_1=0}^{n-1} \delta_{n_1}^{(0)} \right. \\ & + c_2 \sum_{n_1=0}^{n-1} \delta_{n_1}^{(0)} \sum_{n_2=0}^{n_1-1} \delta_{n_2}^{(1)} \\ & \left. + \dots \right)\end{aligned}$$

at level k .

FORMAL DESCRIPTION

- Simple model $\{\sigma_n(\vec{c}, \vec{p})\} \rightarrow \sigma(\vec{p})$
- Simple transformation

$$T(\vec{p}) : \{\sigma_n(\vec{c}, \vec{p})\}_{n=0}^{\infty} \longrightarrow \{\sigma(\vec{p})\}_{n=0}^{\infty}$$

- Hierarchy of model sequences

$$\{\{\sigma_n^{(\ell)}(\vec{c}^{(\ell)}, \vec{p}^{(\ell)}) \mid \vec{c}^{(\ell)} \in \mathbb{C}^{a^{(\ell)}}\}_{n=0}^{\infty}\}_{\ell=0}^L$$

with $a^{(\ell)} > a^{(\ell')}$ for $\ell > \ell'$

- Mapping between levels

$$\begin{aligned} T(\vec{p}^{(\ell)}) : \{\sigma_n^{(\ell)}(\vec{c}^{(\ell)}, \vec{p}^{(\ell)})\}_{n=0}^{\infty} \\ \longrightarrow \{\sigma_n^{(\ell-1)}(\vec{c}^{(\ell-1)}, \vec{p}^{(\ell-1)})\}_{n=0}^{\infty} \end{aligned}$$

- Hierarchical consistent transformation

$$T^{(\ell)} = T(\vec{p}^{(0)}) \circ T(\vec{p}^{(1)}) \circ \dots \circ T(\vec{p}^{(\ell)})$$

METHODS FOR ORTHOGONAL SERIES

- \mathcal{I} transformation:

- Simple model sequence:

$$\sigma_n = \sigma + \omega_n (c \exp(in\nu) + d \exp(-in\nu))$$

- Simple sequence transformation

$$s'_n = \frac{\frac{s_{n+2}}{\omega_{n+2}} - 2 \cos(\nu) \frac{s_{n+1}}{\omega_{n+1}} + \frac{s_n}{\omega_n}}{\frac{1}{\omega_{n+2}} - 2 \cos(\nu) \frac{1}{\omega_{n+1}} + \frac{1}{\omega_n}}$$

- Compute ω'_n ?

- More complicated model sequence:

$$\sigma_n = \sigma + \omega_n (e^{in\nu} (c_0 + c_1 r_n) + e^{-in\nu} (d_0 + d_1 r_n))$$

yields

$$\sigma'_n \approx \sigma + \omega'_n (c' \exp(in\nu) + d' \exp(-in\nu))$$

$$\omega'_n = \frac{-(r_{n+1} - r_n)}{\frac{1}{\omega_{n+2}} - 2 \cos(\nu) \frac{1}{\omega_{n+1}} + \frac{1}{\omega_n}}$$

– Iteration $\implies \mathcal{I}$ transformation

$$s_n^{(0)} = s_n, \quad \omega_n^{(0)} = \omega_n$$

$$s_n^{(k+1)} = \frac{\frac{s_{n+2}^{(k)}}{\omega_{n+2}^{(k)}} - 2 \cos(\nu) \frac{s_{n+1}^{(k)}}{\omega_{n+1}^{(k)}} + \frac{s_n^{(k)}}{\omega_n^{(k)}}}{\frac{1}{\omega_{n+2}^{(k)}} - 2 \cos(\nu) \frac{1}{\omega_{n+1}^{(k)}} + \frac{1}{\omega_n^{(k)}}}$$

$$\omega_n^{(k+1)} = \frac{-\Delta r_{n+1}^{(k)}}{\frac{1}{\omega_{n+2}^{(k)}} - 2 \cos(\nu) \frac{1}{\omega_{n+1}^{(k)}} + \frac{1}{\omega_n^{(k)}}}$$

– for Fourier series

– Notice three-term recurrence

$$u_{n+2} - 2 \cos(\nu) u_{n+1} + u_n = 0$$

satisfied by $\exp(\pm in\nu)$

(or $\cos(n\nu)$, $\sin(n\nu)$)

• \mathcal{K} transformation:

– Simple model sequence :

$$\sigma_n = \sigma + c\omega_n P_n(\cos(\nu))$$

– Three-term recurrence

$$\zeta_n^{(0)} P_n + \zeta_n^{(1)} P_{n+1} + \zeta_n^{(2)} P_{n+2} = 0 .$$

(ν dependent)

– Algorithm (analog to \mathcal{I} transformation)

$$\begin{aligned} s_n^{(0)} &= s_n , & \omega_n^{(0)} &= \omega_n , \\ s_n^{(k+1)} &= \frac{\zeta_{n+k}^{(0)} \frac{s_n^{(k)}}{\omega_n^{(k)}} + \zeta_{n+k}^{(1)} \frac{s_{n+1}^{(k)}}{\omega_{n+1}^{(k)}} + \zeta_{n+k}^{(2)} \frac{s_{n+1}^{(k)}}{\omega_{n+1}^{(k)}}}{\zeta_{n+k}^{(0)} \frac{1}{\omega_n^{(k)}} + \zeta_{n+k}^{(1)} \frac{1}{\omega_{n+1}^{(k)}} + \zeta_{n+k}^{(2)} \frac{1}{\omega_{n+1}^{(k)}}} \\ \omega_n^{(k+1)} &= \frac{\delta_n^{(k)}}{\zeta_{n+k}^{(0)} \frac{1}{\omega_n^{(k)}} + \zeta_{n+k}^{(1)} \frac{1}{\omega_{n+1}^{(k)}} + \zeta_{n+k}^{(2)} \frac{1}{\omega_{n+1}^{(k)}}} \\ \mathcal{K}_n^{(k)}(\{\delta_n^{(k)}\}, \{\zeta_n^{(j)}\}, \{s_n\}, \{\omega_n\}) &= s_n^{(k)} \end{aligned}$$

– ν dependent, for orthogonal series.

MULTIPOLE EXPANSIONS

$$U_Q(\vec{r}) = 4\pi \sum_{\ell m} \frac{1}{r^{\ell+1}} \frac{Y_\ell^m(\vec{r}/r)}{2\ell+1} Q_\ell^{m*}$$

$$Q_\ell^m = \int r'^\ell Y_\ell^m(\vec{r}'/r') \rho(\vec{r}') d^3r',$$

Rotational symmetry

$$U_Q(\vec{r}) = \sum_{\ell=0}^{\infty} P_\ell \left(\frac{\vec{r} \cdot \vec{R}}{r R} \right) \frac{q_\ell}{r^{\ell+1}}$$

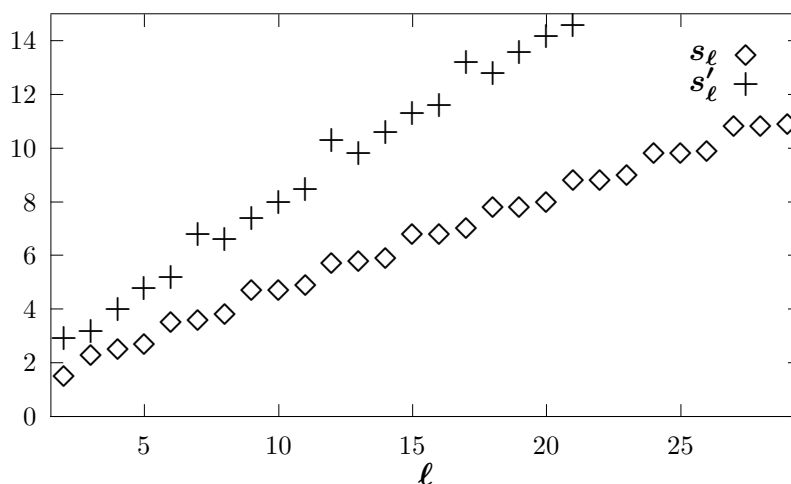
Legendre Expansion

Example

$$\rho(\vec{r}) = \exp(-\alpha r) \exp(-\beta |\vec{r} - \vec{R}|)$$

ℓ	$-\lg 1 - s_\ell/s $	$-\lg 1 - s'_\ell/s $
2	2.6	5.1
4	4.3	9.7
6	5.9	11.5
8	7.6	16.0
10	9.2	16.0

\mathcal{K} transformation, $r = 12$, $\theta = 60^\circ$



Exact digits ($r = 4$, $\theta = 60^\circ$)

MANYFOLD FREQUENCIES

- Important near singularities
- Increases stability
- Instead of

$$s_0, s_1, s_2, \dots, \omega_0, \omega_1, \omega_2, \dots$$

take

$$s_{\tau \cdot 0}, s_{\tau \cdot 1}, s_{\tau \cdot 2}, \dots, \omega_{\tau \cdot 0}, \omega_{\tau \cdot 1}, \omega_{\tau \cdot 2}, \dots$$

- for Fourier and orthogonal series put

$$\nu \rightarrow \tau \cdot \nu$$

i.e., for $x = \cos \nu$

$$x \rightarrow x_\tau = \cos(\tau \cdot \arccos x)$$

$$\sum_{\ell=0}^{\infty} \frac{1}{\ell+1} P_{\ell}(x) = \ln \left(1 + \sqrt{\frac{2}{1-x}} \right)$$

$$x = 0.9$$

\mathcal{K} transformation

$$\tau = 1$$

n	$-\lg 1 - s_n/s $	$-\lg 1 - s'_n/s $
16	2.07	5.24
18	1.75	6.88
20	1.91	6.58
22	3.59	6.91
24	2.01	6.80

$$\tau = 3$$

m	n	$-\lg 1 - s_m/s $	$-\lg 1 - s'_n/s $
48	16	2.51	9.31
54	18	2.45	10.40
60	20	2.48	11.63
66	22	2.59	13.18
72	24	2.80	14.47

ASSOCIATED POWER SERIES

Example:

$$\sum_{n=0}^{\infty} \frac{1+in}{n^2} \cos(n\nu) = \frac{1}{2} \sum_{n=0}^{\infty} \frac{1+in}{n^2} e^{in\nu} + \frac{1}{2} \sum_{n=0}^{\infty} \frac{1+in}{n^2} e^{-in\nu}$$

Accelerate power series separately and add

Adaptable to more complicated examples:

$$\sum \cos((n+1/2)\nu) P_n(\cos \nu')$$

(singular at $\nu = \nu'$) is sum of 4 power series

$$\frac{1}{4} \sum_{n=0}^{\infty} e^{\pm i(n+1/2)\nu} \rho_n^{\pm}(\nu')$$

$$\rho_n^{\pm}(\nu') = P_n(\cos \nu') \pm i \frac{2}{\pi} Q_n(\cos \nu') \\ \sim \frac{\exp(\pm in\nu')}{\sqrt{n}} \times \text{const.}$$

$$\tau = 10, \quad \nu = 6\pi/10, \quad \nu' = 2\pi/3$$

near singularity.

n	$-\lg (s_{\tau n} - s)/s $	$-\lg (G_n^{(\tau)} - s)/s $
8	1.3	7.7
12	1.2	14.0
16	1.0	18.1
20	1.5	22.6
24	1.3	27.3
28	1.2	31.2

$$G_n^{(\tau)} = \sum_{j=1}^4 \mathcal{L}_n^{(0)}(1, [p_{j,\tau n}]|_{n=0}, [(\tau n + 1)(p_{j,\tau n} - p_{j,(\tau n)-1})]|_{n=0})$$

VARIATIONAL METHODS

Problem:

Exact limit invariant under addition of null sequences

Nonlinear sequence transformation usually not !

Idea:

Restore invariance variationally for certain null sequences x_n

$$s'_n = f(\{s_n + \alpha x_n\})$$

$$\frac{\partial s'_n}{\partial \alpha} = 0$$

Example:

Aitken Δ^2 method

$$s_n^{(1)} = s_n - \frac{(\Delta s_n)^2}{\Delta^2 s_n}.$$

Use $s_n \rightarrow s_n + \alpha x_n$ with $\lim_{n \rightarrow \infty} x_n = 0$

$$s_n^{(1)} \rightarrow s_n^{(1)}(\alpha) = s_n + \alpha x_n - \frac{(\Delta s_n + \alpha \Delta x_n)^2}{\Delta^2 s_n + \alpha \Delta x_n}.$$

Choose α such that $s_n^{(1)}(\alpha)$ is stationary:

$$\left. \frac{\partial s_n^{(1)}}{\partial \alpha} \right|_{\alpha=\alpha_0} = 0$$

Result: For $s_n = 10 + 1/n^2$ new method accelerates ($O(n^{-3})$ error), but Aitken does not(!) accelerate convergence.

Example

$$s'_n = \sum_{j=0}^k c_j s_{n+j}, \quad \sum_{j=0}^k c_j = 1$$

Put

$$s_n \rightarrow s_n + \sum_{\nu=1}^k \alpha_\nu x_n^{(\nu)}$$

Saddle point \implies Linear system

$$\frac{\partial s'_n}{\partial \alpha_\mu} = \sum_{j=0}^k c_j x_{n+j}^{(\mu)} = 0$$

Result identical to E algorithm !

$$s'_n = \frac{\begin{vmatrix} s_n & \cdots & s_{n+k} \\ x_n^{(1)} & \cdots & x_{n+k}^{(1)} \\ \vdots & \cdots & \vdots \\ x_n^{(k)} & \cdots & x_{n+k}^{(k)} \end{vmatrix}}{\begin{vmatrix} 1 & \cdots & 1 \\ x_n^{(1)} & \cdots & x_{n+k}^{(1)} \\ \vdots & \cdots & \vdots \\ x_n^{(k)} & \cdots & x_{n+k}^{(k)} \end{vmatrix}}$$

PERTURBATIONAL METHODS

Rayleigh-Schrödinger Perturbation Theory

$$H = H_0 + \beta V$$

yields

$$E^{(n)} = E_0 + \beta E_1 + \dots + \beta^n E_n$$

Goldhammer-Feenberg

$$H = (1 - \alpha)H_0 + [V + \alpha H_0]$$

yields

$$E^{(n)}(\alpha) = E_0(\alpha) + \beta E_1(\alpha) + \dots + \beta^n E_n(\alpha)$$

Choose α variationally

(True E is α -independent)

$$\frac{\partial E^{(n)}(\alpha)}{\partial \alpha} = 0$$

For $n = 3$ solution is $\alpha = E_3/E_2 \longrightarrow$ Feenberg series

$$F_n = E^{(n)}(E_3/E_2)$$

EFFECTIVE CHARACTERISTIC POLYNOMIALS

Čížek

$$\begin{aligned}
 P_n(E) &= \det \left| \langle \phi_j | H | \phi_k \rangle - E \delta_{j,k} \right| \\
 &= \sum_{j=0}^n E^j \sum_{k=0}^{n-j} f_{n,j,k} \beta^k
 \end{aligned}$$

Obtain f 's from perturbation series

$$P_n(E_0 + \beta E_1 + \beta^2 E_2 + \dots) = O(\beta^{n(n+3)/2})$$

Zero of P_2 :

$$\begin{aligned}
 \Pi_2 &= E_0 + E_1 \\
 &+ \frac{E_2^2}{2} \frac{E_2 - E_3}{E_2 E_4 - E_3^2} \\
 &+ \frac{E_2^2}{2} \frac{\sqrt{(E_2 - E_3)^2 - 4(E_2 E_4 - E_3^2)}}{E_2 E_4 - E_3^2}
 \end{aligned}$$

Invariant under Feenberg scaling

$$\Pi_2(E_0, \dots, E_4) = \Pi_2(E_0(\alpha), \dots, E_4(\alpha)) .$$

Scaling property

$$\Pi_2(c E_0, \dots, c E_4) = c \Pi_2(E_0, \dots, E_4) .$$

MANY-BODY PERTURBATION THEORY

Dissociation barrier (kJ/mol) for
 $\text{H}_2\text{CO} \longrightarrow \text{H}_2 + \text{CO}$

Method	Minimum	Transition state	Barr.
$E_0 + E_1$	-113.912879	-113.748693	431.1
E_2	-114.329202	-114.182435	385.3
E_3	-114.334186	-114.185375	390.7
E_4	-114.359894	-114.219892	367.6
F_4	-114.360838	-114.220603	368.2
[2/2]	-114.362267	-114.223409	364.6
Π_2	-114.364840	-114.227767	359.9
Best Estimate			360

(TZ2P Basis at MP2 Geometries)

ITERATION SEQUENCES

Fixed-point equation

$$x = \Psi(x)$$

Direct Iteration

$$x_0, x_1 = \Psi(x_0), \dots, x_{n+1} = \Psi(x_n), \dots$$

Cycling

$$s_0 = x_{\text{Start}}, s_1 = \Psi(s_0), \dots, s_k = \Psi(s_{k-1})$$
$$x_{\text{Start}} = \mathcal{T}(s_0, \dots, s_k)$$

Corresponds to new iteration function:

$$y_{n+1} = \mathcal{T}(y_n, \Psi(y_n), \dots, \Psi(\Psi(\dots \Psi(y_n))))$$

ORNSTEIN-ZERNIKE-EQUATION

- Classical many-particle systems (fluids)
- Pair distribution function $g(\mathbf{r}) = 1 + h(\mathbf{r})$

- Integral equation

$$h = c + \rho h * c$$

$$g(\mathbf{r}) = \exp(-\beta u(\mathbf{r}) + h(\mathbf{r}) - c(\mathbf{r}) + E(\mathbf{r}))$$

- Bridge diagrams $E(\mathbf{r}) \implies$ Closure relations

- Solution on lattice with FFT:
 $\Gamma_i = (h(r_i) - c(r_i))r_i, r_i = i\Delta r$

$$\vec{\Gamma} = \Psi(\vec{\Gamma})$$

via

- direct iteration
- direct iteration
- + vector extrapolation
- Extrapolation reduces CPU time by up to 50 %
- Extrapolation useful to achieve convergence

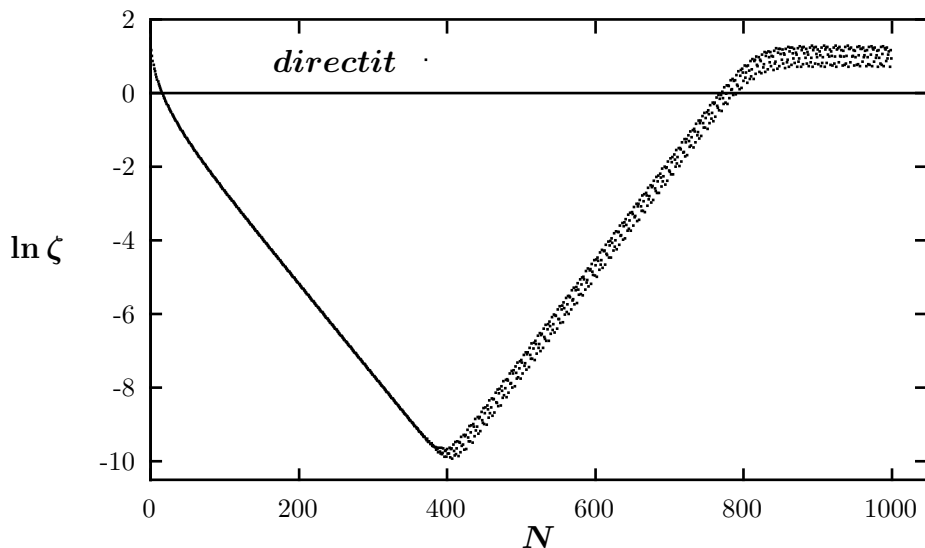


Figure 1: Unstable Fixed-point of Direct Iteration (Hard spheres, high density)

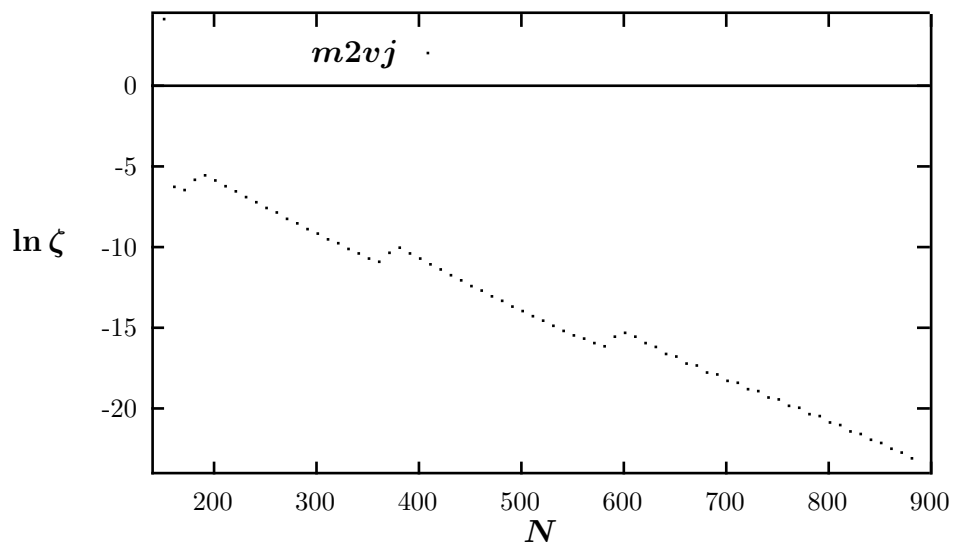


Figure 2: Convergence of the Cyclig Method