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Transverse momentum dependent quark densities from Lattice QCD

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Abstract. We study transverse momentum dependent parton distribution functions (TMDs) with non-local operators in lattice QCD, using MILC/LHPC lattices. We discuss the basic concepts of the method, including renormalization of the gauge link. Results obtained with a simplified operator geometry show visible dipole deformations of spin-dependent quark momentum densities.

Keywords: transverse momentum; parton distribution functions; lattice; QCD
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INTRODUCTION

Generalized parton distribution functions (GPDs) and transverse momentum dependent parton distribution functions (TMDs) provide us with a picture of the internal quark distributions in a nucleon at the instant of an interaction, see illustration Fig. 1 a). GPDs and TMDs have their natural interpretation at large nucleon momentum $\mathbf{P} = (0,0,\mathbf{P}_z)$. The quark momentum k in terms of light cone coordinates $k^{\pm} \equiv (k^0 \pm k^3)/\sqrt{2}$, $\mathbf{k}_\perp = (\mathbf{k}_x, \mathbf{k}_y)$ scales like $k^+ : \mathbf{k}_\perp : k^- \sim P^+ : 1 : (P^+)^{-1}$ with the large momentum component P^+ of the nucleon. TMDs resolve the dependence on $x \equiv k^+/P^+$ and transverse momentum $\mathbf{k}_\perp$, but not on the suppressed component k^-. In spin-polarized channels at leading twist, TMDs encode dipole- or quadrupole-shaped deformations of the nucleon in the $\mathbf{k}_\perp$-plane. We have studied such deformations in first explorative lattice QCD calculations [1, 2, 3], see Fig. 1 and our discussion below. These studies have been motivated by a history of successful lattice computations of x-moments of GPDs, providing images of the nucleon in the impact parameter, $\mathbf{b}_\perp$-, plane, see [4] for a review. A remaining theoretical problem concerns the precise form of the correlator defining TMDs in the continuum, see [5, 6] and references therein. In its basic form, it is given by [7]

$$\Phi_q^{[\Gamma]}(x,\mathbf{k}_\perp;P,S;\mathscr{C}) \equiv \int dk^- \int \frac{d^4 l}{(2\pi)^4}\, e^{-ik\cdot l}\, \underbrace{\frac{1}{2}\langle P,S|\ \bar{q}(l)\,\Gamma\,\mathscr{U}[\mathscr{C}_l]\ q(0)\ |P,S\rangle}_{\widetilde{\Phi}_q^{[\Gamma]}(l,P,S;\mathscr{C})}\Big|_{k^+=xP^+}$$
$$= \frac{1}{P^+}\underbrace{\int \frac{d(l\cdot P)}{2\pi}\, e^{-i(l\cdot P)x}}_{\not{\int}_x}\ \underbrace{\int \frac{d^2\mathbf{l}_\perp}{(2\pi)^2}\, e^{i\mathbf{l}_\perp\cdot\mathbf{k}_\perp}}_{\not{\int}_{\mathbf{l}}}\widetilde{\Phi}_q^{[\Gamma]}(l,P,S;\mathscr{C})\Big|_{l^+=0} \qquad (1)$$

where Γ is a Dirac matrix. The Wilson line $\mathscr{U}[\mathscr{C}_l]$ running along a continuous path $\mathscr{C}_l$ from l to 0 ensures gauge invariance of the expression. For the SIDIS and Drell-Yan scattering process, the Wilson line extends to infinity along a direction v that needs to be chosen (almost) lightlike, such that the cross section factorizes into hard, perturbative parts and soft contributions, see, e.g., Ref. [8]. Based on its symmetry transformation properties, the above correlator can be parametrized in terms of TMDs [9, 10, 11], for example

$$2\,\rho_{TL}^{(q)} \equiv \Phi_q^{[\gamma^+ + \lambda\gamma^+\gamma^5]} = f_{1,q} + \lambda\,\frac{\mathbf{k}_\perp\cdot\mathbf{S}_\perp}{m_N} g_{1T,q} + \left[\frac{\mathbf{S}_j\varepsilon_{ji}\mathbf{k}_i}{m_N} f_{1T,q}^{\perp}\right]_{\text{odd}}, \qquad (2)$$

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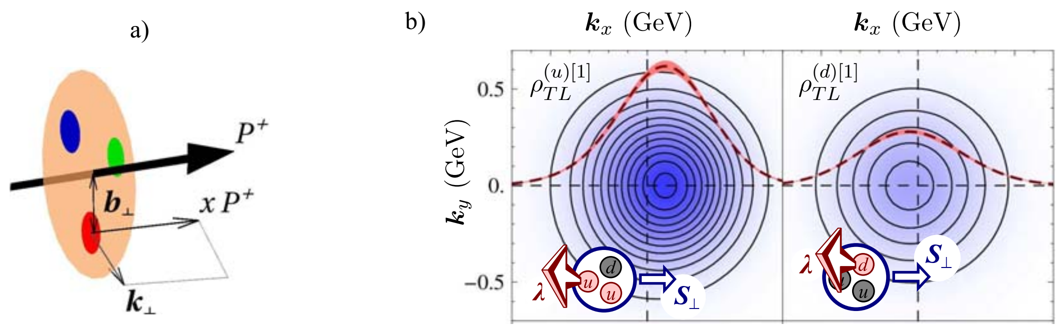

FIGURE 1. a) Illustration of quark degrees of freedom in the nucleon at large momentum. b) Dipole-deformed x-integrated densities obtained with straight gauge links at a pion mass $m_\pi \approx 500\,\mathrm{MeV}$. The insets display the spin polarization of the quarks (red arrow) and of the nucleon (blue arrow).

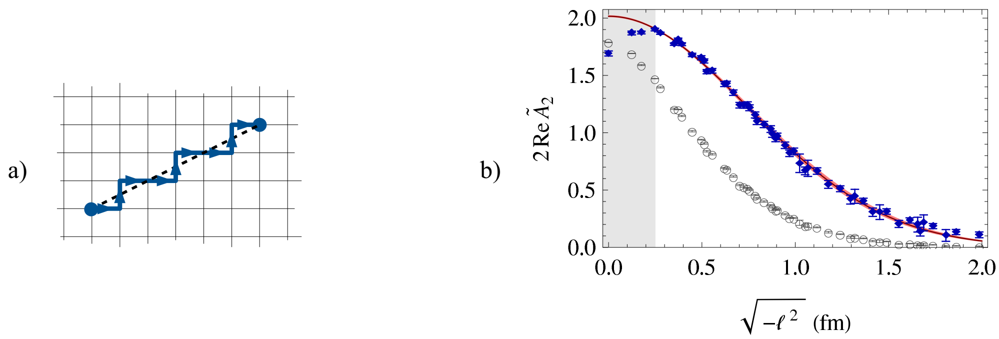

FIGURE 2. a) Representation of a straight Wilson line (dashed line) as a step-like product of link variables. b) Amplitude $\widetilde{A}_2(l^2,0)$ for up quarks at a pion mass $m_\pi \approx 500\,\mathrm{MeV}$, using straight gauge links.

Here λ is the longitudinal quark polarization, and Λ and $\mathbf{S}_\perp$ are longitudinal and transverse nucleon polarization, respecively. The leading-twist TMDs $f_{1,q}, g_{1T,q}, f_{1T,q}^\perp$ are real-valued functions of x and $\mathbf{k}_\perp^2$. The "naively time-reversal odd" function $f_{1T,q}^\perp$ switches sign when comparing the SIDIS- with the Drell-Yan process, because the direction v of the Wilson line changes from future- to past-pointing [12].

STRAIGHT LINK TMDS FROM THE LATTICE

In light of the uncertainties about the precise form of the continuum correlator, and to develop our methods, our first lattice studies employ a simple operator geometry that does not relate to a specific scattering process: We connect the quark fields with a direct, straight Wilson line. For the resulting "process-independent" TMDs, the T-odd functions such as the Sivers function $f_{1T,q}^\perp$ vanish exactly.

In our approach, we calculate matrix elements $\langle P,S|\,O\,|P,S\rangle$ from ratios of three- and two-point functions using the same techniques as GPD calculations by the LHP collaboration in Ref. [13]. We also use the same sequential propagators and quark propagators, calculated by LHPC with domain-wall valence fermions on top of asqtad-improved staggered MILC gauge configurations [14, 15, 16] with 2+1 quark flavors at a lattice spacing $a \approx 0.12\,\mathrm{fm}$. The difference with respect to GPD calculations is that we directly insert the non-local operator $O \equiv \bar{q}(l)\Gamma\mathscr{U}[\mathscr{C}_l]q(0)$ in our three-point function. The Wilson line $\mathscr{U}[\mathscr{C}_l]$ is approximated as a step-like product of HYP-smeared link-variables as illustrated in Fig. 2 a). See also Ref. [2, 3].

The connection between the matrix elements $\tilde{\Phi}^{[\Gamma]}$ and TMDs is established through a parametrization in terms of Lorentz-invariant amplitudes $\widetilde{A}_i(l^2, l\cdot P)$. For straight Wilson lines, we obtain in analogy to the parametrization in terms

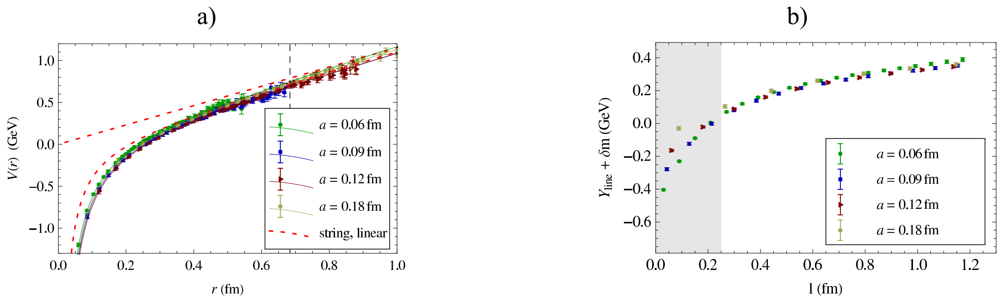

FIGURE 3. a) Static quark potential from MILC lattices at several lattice spacings a, matched to the string potential at $r \approx 0.7$ fm. b) Test of the renormalization procedure with straight Wilson lines on a gauge fixed ensemble.

of amplitudes $A_i(k^2, k\cdot P)$ in Ref. [9] (here our sign conventions follow Ref. [11] with the substitution rule $k \to i m_N^2 l$):

$$\tilde{\Phi}^{[\gamma^\mu]} = 2P^\mu \widetilde{A}_2 + 2i m_N{}^2 l^\mu \widetilde{A}_3\,, \qquad \tilde{\Phi}^{[\gamma^\mu\gamma^5]} = -2m_N S^\mu \widetilde{A}_6 - 2i m_N P^\mu (l\cdot S)\widetilde{A}_7 + 2m_N{}^3 l^\mu (l\cdot S)\widetilde{A}_8\,.$$

The TMDs are then obtained by

$$f_1(x,\mathbf{k}_\perp^2) = 2 \not{\int}_{\!X} \not{\int}_{\!M} \widetilde{A}_2(l^2, l\cdot P)\,, \qquad g_{1T}(x,\mathbf{k}_\perp^2) = 4m_N^2 \partial_{\mathbf{k}_\perp^2} \not{\int}_{\!X} \not{\int}_{\!M} \widetilde{A}_7(l^2, l\cdot P)\,.$$

In the equations above, $\not{\int}_{\!X}$ only acts on $l\cdot P$, while $\not{\int}_{\!M}$ only acts on l^2. Thus $x \leftrightarrow l\cdot P$ and $\mathbf{k}_\perp^2 \leftrightarrow l^2$ are pairs of conjugate variables. Our Euclidean lattice approach is restricted to the determination of amplitudes $\widetilde{A}_i$ for $l^0 = -il_4 = 0$, i.e., to the region $l^2 < 0$, $|l\cdot P| \leq \sqrt{-l^2}|\mathbf{P}|$, where $\mathbf{P}$ is the selected three-momentum of the nucleon on the lattice. The limited range in $|l\cdot P|$ prohibits us from a direct evaluation of $\not{\int}_{\!X}$. However, first studies of x- and $\mathbf{k}_\perp$- correlations are possible [17, 3]. Moreover, x-integrated TMDs and densities are directly accessible: Integrating Eq. (1) with respect to x removes $\not{\int}_{\!X}$ and sets $l\cdot P$ to zero. Correspondingly, the x-integral of, e.g., f_1 becomes $\int_{-1}^{1} dx\, f_1(x,\mathbf{k}_\perp^2) \equiv f_1^{[1]}(\mathbf{k}_\perp^2) = 2\not{\int}_{\!M} \widetilde{A}_2(l^2,0)$. In Fig. 2 b), open symbols correspond to unrenormalized lattice data for $\widetilde{A}_2(l^2,0)$.

To obtain results independent of our lattice spacing a and our lattice action, we must renormalize our data. The Wilson line $\mathscr{U}[\mathscr{C}_l]$ introduces a length dependent renormalization factor $\exp(-\delta m\sqrt{-l^2})$ [18, 19, 20]. To fix δm, we follow the strategy of Refs. [21, 22], and match the renormalized static quark potential $V^{\text{ren}}(r) = V(r) + 2\delta m$ to the string potential $V_{\text{string}} = \sigma r - \pi/(12r)$ [23] at a matching point $r = 1.5 r_0 \approx 0.7$ fm. In Fig. 3 a), we test the method for several lattice spacings a on four MILC lattices with similar pion masses $m_\pi \approx 500$ MeV. The renormalized lattice data agree very well with each other and are approximated well by the string potential (red dashed curve) near the matching point, indicated by a vertical dashed line. The procedure implements a gauge-invariant renormalization condition that we can formulate as the demand that the static quark potential asymptotically approach a straight line σr through the origin (shown as a red dashed line). In connection with TMDs, we lack at present an interpretation of this renormalization condition as a physical renormalization or factorization scale. In Figure 3 b), we check the applicability of the approach to Wilson lines by plotting $Y_{\text{line}}^{\text{ren}}(l) = \ln(U_{l-a/2}/U_{l+a/2})/a + \delta m$, where U_l is the expectation value of the color trace of a straight Wilson line of length l evaluated on a Landau gauge fixed ensemble, and where the length dependent renormalization has been carried out with the values δm obtained from the static quark potential. Only at short lengths, $l \lesssim 0.25$ fm, we find significant differences between lattice data from different lattice spacings, a sign of lattice cutoff effects. For our TMD calculations discussed below we exclude data obtained in this region from our fits. For $l \gtrsim 0.25$ fm, we assume that renormalization of the lattice operator can be carried out as in the continuum, $O^{\text{ren}} = Z_{\Psi,z}^{-1} \exp(-\delta m\sqrt{-l^2})\, O$, where the renormalization constants $Z_{\Psi,z}^{-1}$ and δm are independent of the Dirac structure Γ [19].

Figure 2 b) shows the renormalized lattice data for $\widetilde{A}_2(l^2,0)$ as solid data points. The curve and statistical error band correspond to a Gaussian fit to this data in the range $\sqrt{-l^2} \geq 0.25$ fm. Note that the renormalization constant $Z_{\Psi,z}^{-1}$ has been fixed (in the isovector, $u-d$-channel) such that the x-$\mathbf{k}_\perp$-integrated Gaussian density of unpolarized quarks yields the correct total number of valence quarks, $\int d^2\mathbf{k}_\perp f_{1,u-d}^{[1]} = 1$. Similar fits for $\widetilde{A}_7$ enable us to calculate the "worm-gear" function $g_{1T}^{[1]}$, and correspondingly, the dipole deformed x-integrated density $\rho_{TL}^{(q)[1]}$ defined in Eq. (2) and shown in Fig. 1 b). While the widths of our distributions depend strongly on our renormalization condition for δm, the average transverse quark momentum shift can be expressed in terms of ratios of the Gaussian amplitudes at $l^2=0$:

$$\langle \mathbf{k}_x \rangle_{TL} \equiv \left. \frac{\int d^2\mathbf{k}_\perp \, \mathbf{k}_x \, \rho_{TL}^{[1]}}{\int d^2\mathbf{k}_\perp \rho_{TL}^{[1]}} \right|_{\substack{\lambda=1,\\ \mathbf{S}_\perp=(1,0)}} = m_N \frac{\int d^2\mathbf{k}_\perp \, \mathbf{k}_\perp^2/(2m_N^2)\, g_{1T}^{[1]}(\mathbf{k}_\perp)}{\int d^2\mathbf{k}_\perp \, f_1^{[1]}(\mathbf{k}_\perp)} = -m_N \frac{\widetilde{A}_7(0,0)}{\widetilde{A}_2(0,0)} = \begin{cases} 67(5)\,\mathrm{MeV} & \text{(up)} \\ -30(5)\,\mathrm{MeV} & \text{(down)} \end{cases}$$

(errors statistical only). In these ratios, renormalization factors largely cancel. Reference [24] reveals a remarkable similarity of our results with a light-cone constituent quark model [25], despite the unphysically large quark masses employed in our lattice calculation: They find $\langle \mathbf{k}_x \rangle_{TL} = 55.8\,\mathrm{MeV}$ for up-, and $\langle \mathbf{k}_x \rangle_{TL} = -27.9\,\mathrm{MeV}$ for down-quarks.

CONCLUSIONS AND OUTLOOK

We have performed first lattice studies of TMDs using non-local operators with a simplified, straight gauge link. Resulting average momentum shifts $\langle \mathbf{k}_x \rangle_{TL}$ corroborate model results. An ongoing project with staple-shaped gauge links can potentially address TMDs specific to SIDIS or the Drell-Yan process, including T-odd functions responsible for single-spin asymmetries.

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